

# The $0\nu\beta\beta$ -DGT correlation and wave function factorization

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## Double Gamow-Teller strength distribution

Measurement of Double Gamow-Teller (DGT) resonance  
in double charge-exchange reactions  $^{48}\text{Ca}(pp,nn)^{48}\text{Ti}$  proposed in 80's

Auerbach, Muto, Vogel... 1980's, 90's

Recent experimental plans in RCNP, RIKEN ( $^{48}\text{Ca}$ ), INFN Catania

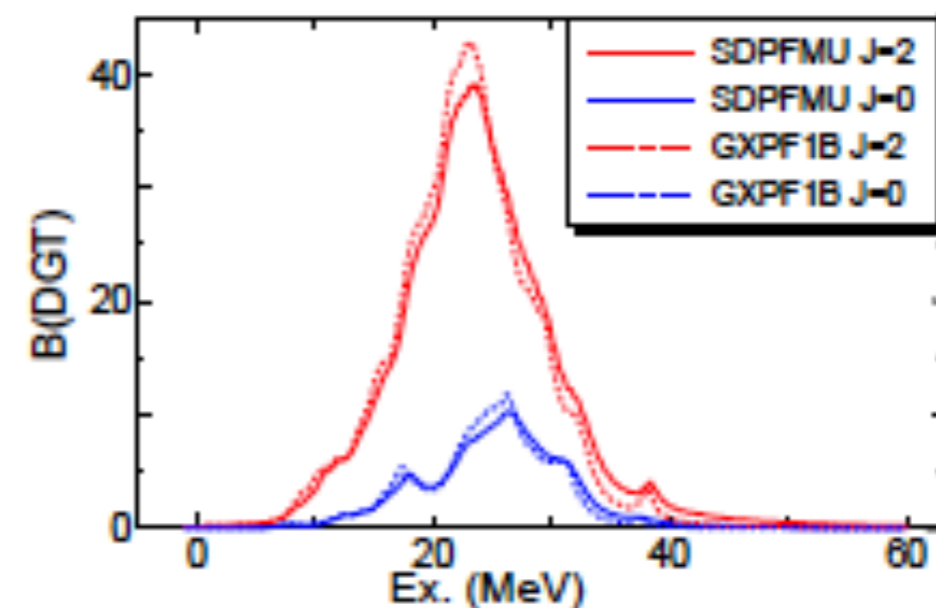
Takaki et al. JPS Conf. Proc. 6 020038 (2015)

Capuzzello et al. EPJA 51 145 (2015), Takahisa, Ejiri et al. arXiv:1703.08264

Promising connection to  $\beta\beta$  decay,  
two-particle-exchange process,  
especially the (tiny) transition  
to ground state of final state

Shell model calculation

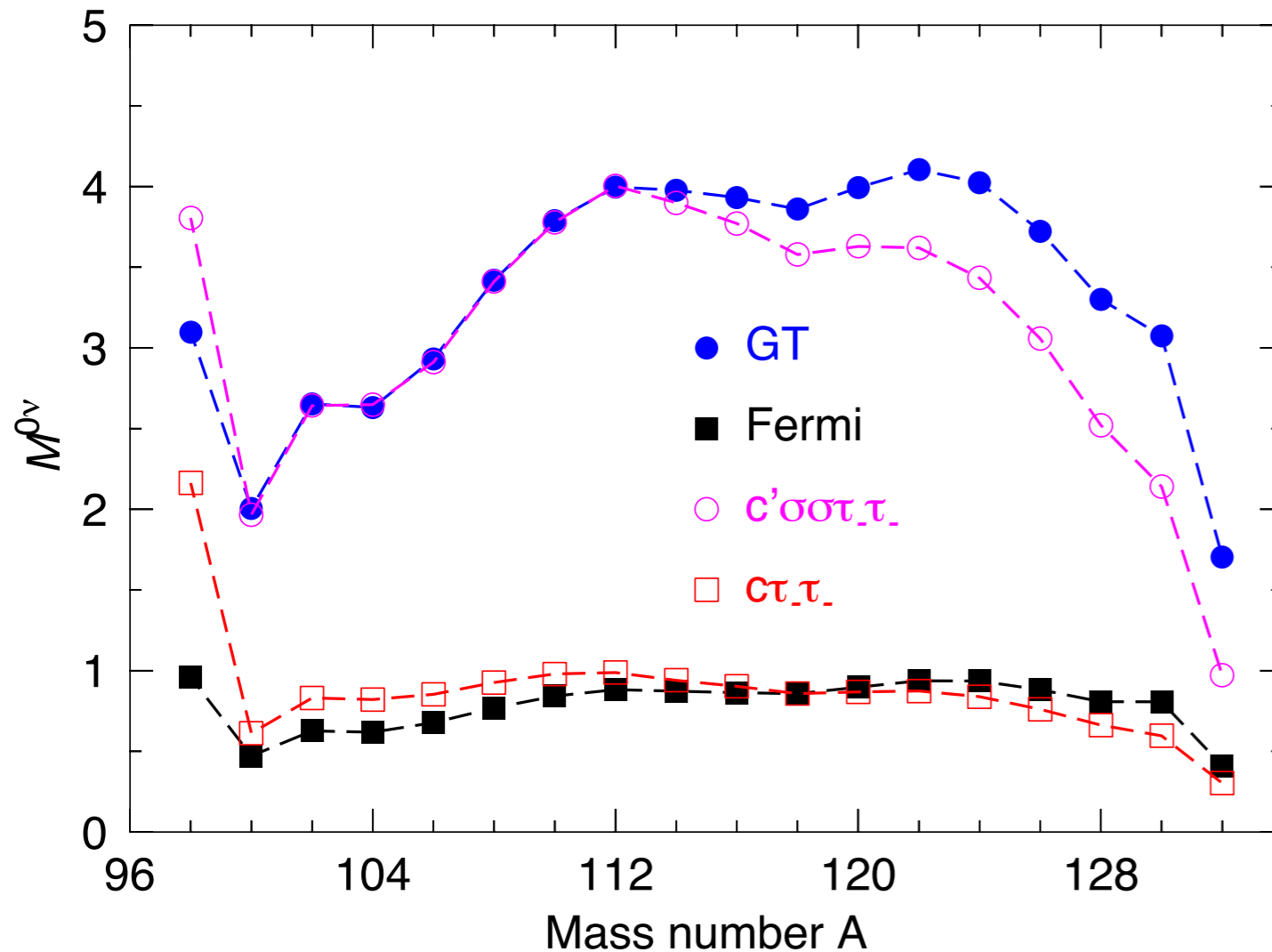
Shimizu, JM, Yako, PRL 120 142502 (2018)



$$B(DGT^-; \lambda; i \rightarrow f) = \frac{1}{2J_i + 1} \left| \left\langle ^{48}\text{Ti} \left\| \left[ \sum_i \sigma_i \tau_i^- \times \sum_j \sigma_j \tau_j^- \right]^{(\lambda)} \right\| ^{48}\text{Ca}_{\text{gs}} \right\rangle \right|^2$$

# $0\nu\beta\beta$ -DGT correlation

T. Rodriguez and G. Martinez-Pinedo PLB 719 (2013)

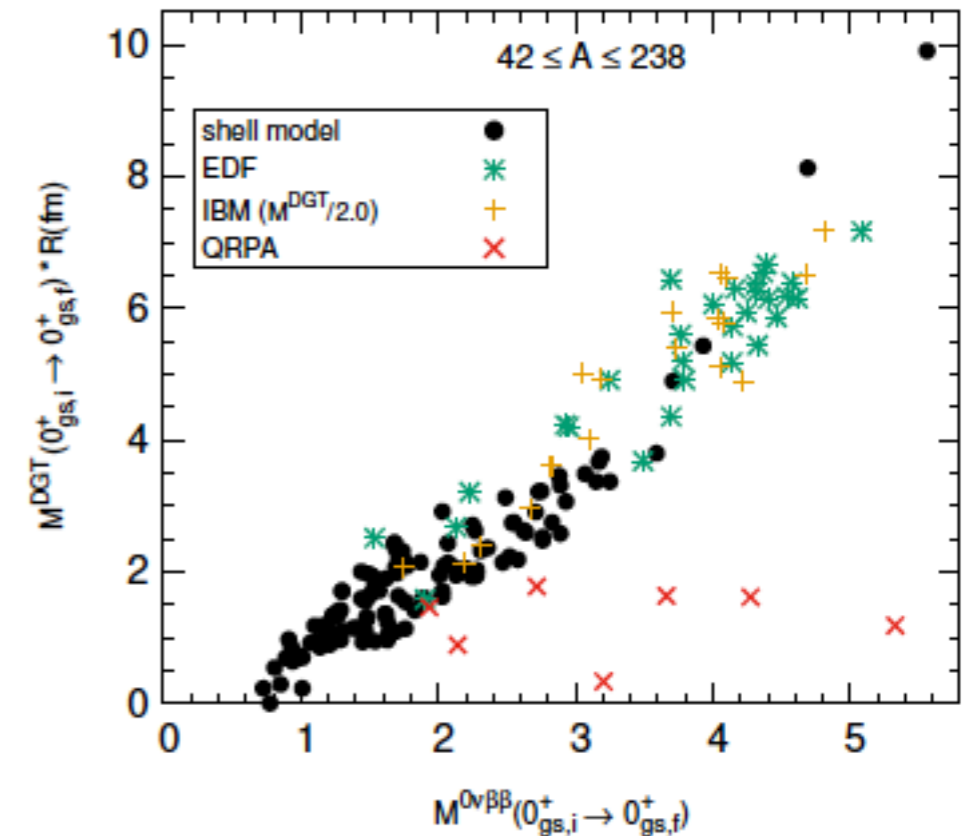
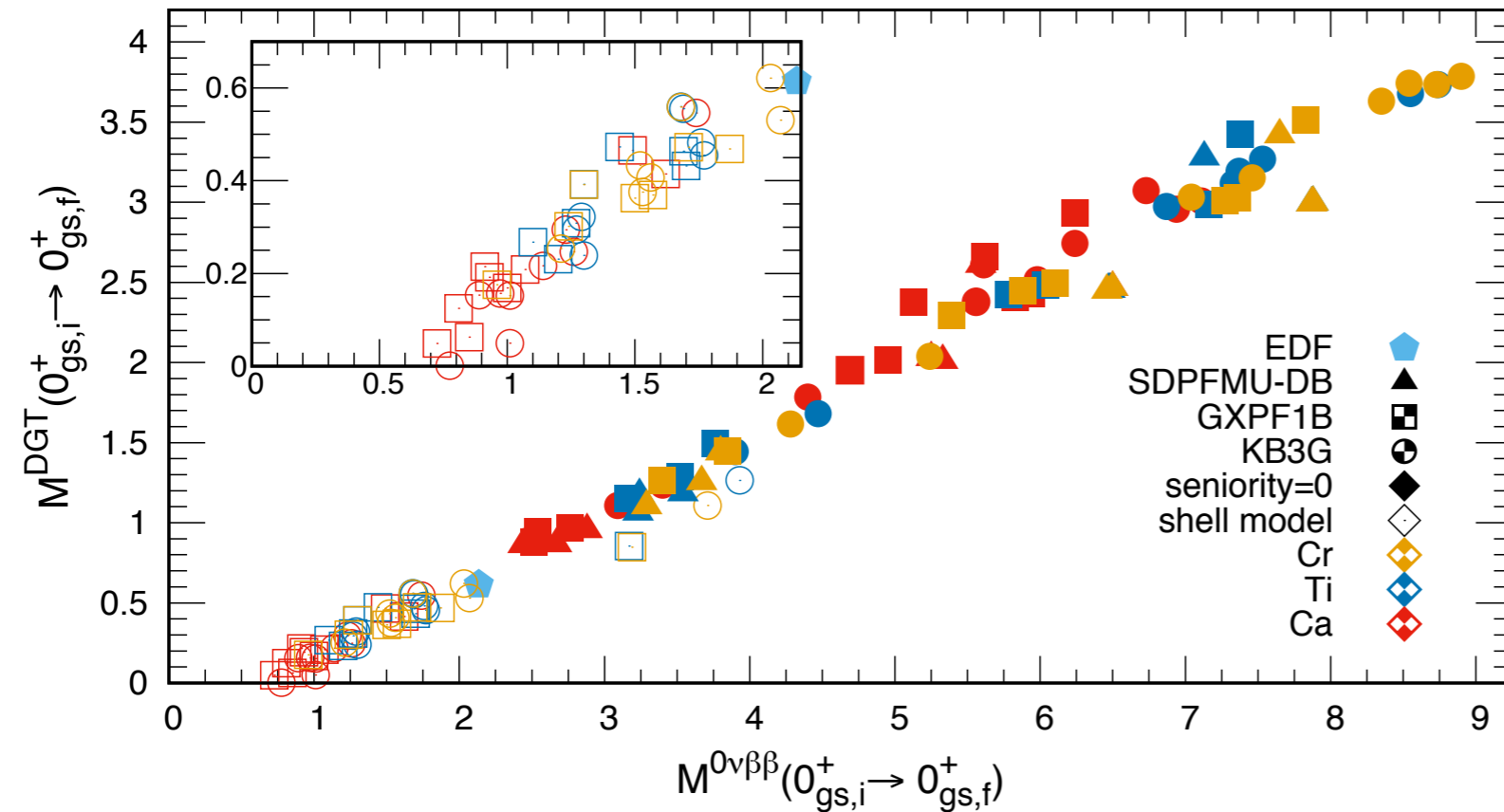


GCM calculations w/Gogny EDF

Match  $0\nu\beta\beta$  matrix elements reasonably well across different  $A$  by replacing  $H_{GT,F}(r)$  with  $c = c' = 2.0$

# $0\nu\beta\beta$ -DGT correlation

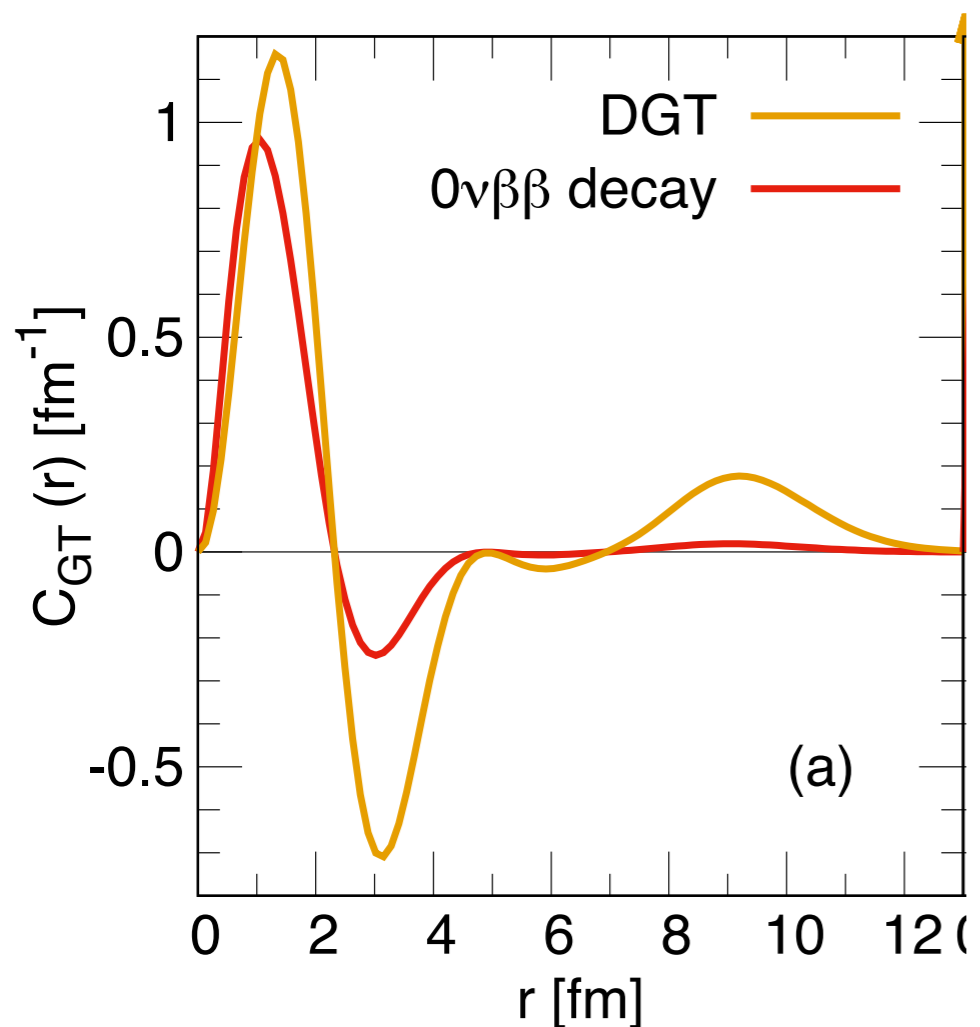
Shimizu, Menendez, Yako PRL 120 (2018)



Correlation robust across the chart ( $42 \leq A \leq 238$ ) for a host of different methods (shell model, IBM, EDF-GCM)

Doesn't hold for QRPA calculations

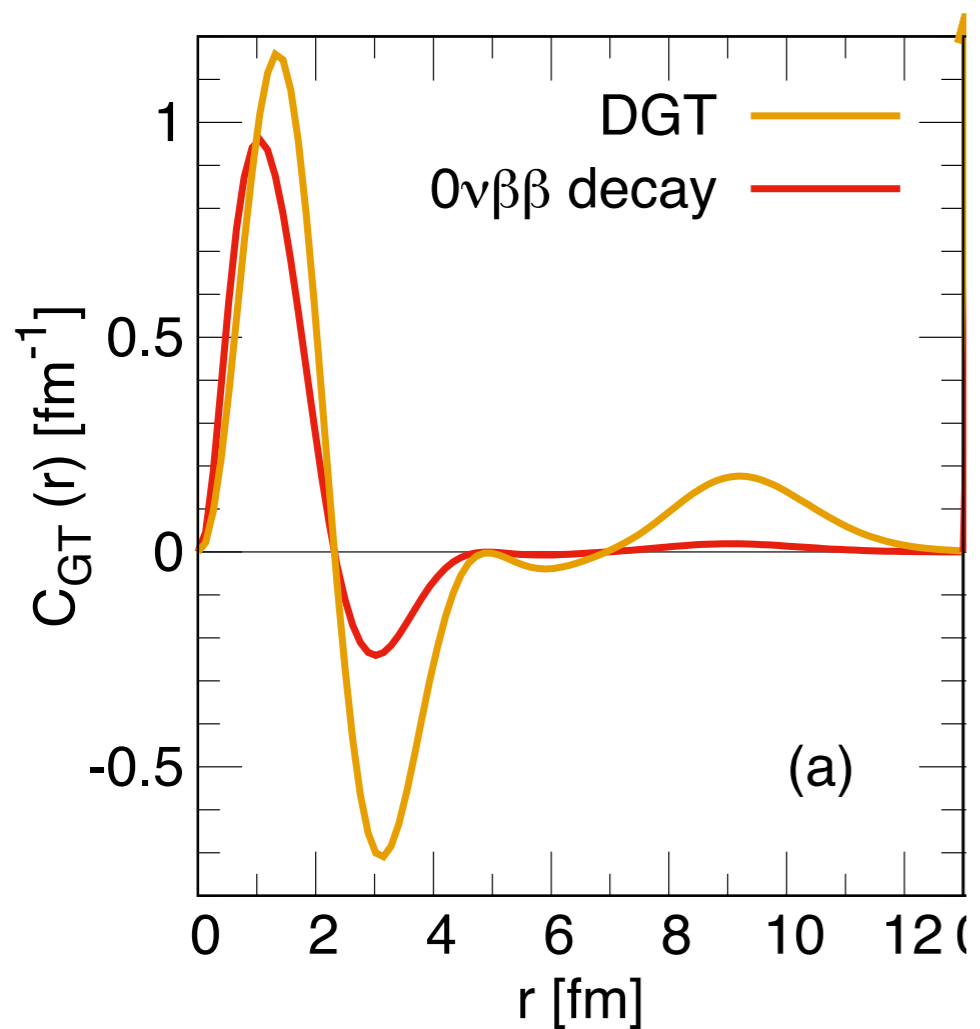
# $0\nu\beta\beta$ -DGT correlation



Dominant contributions to both  $0\nu\beta\beta$  and DGT come from  $r \approx 2$  fm.

Robust cancellation in DGT at large distances, making it “effectively” same range as  $0\nu\beta\beta$ .

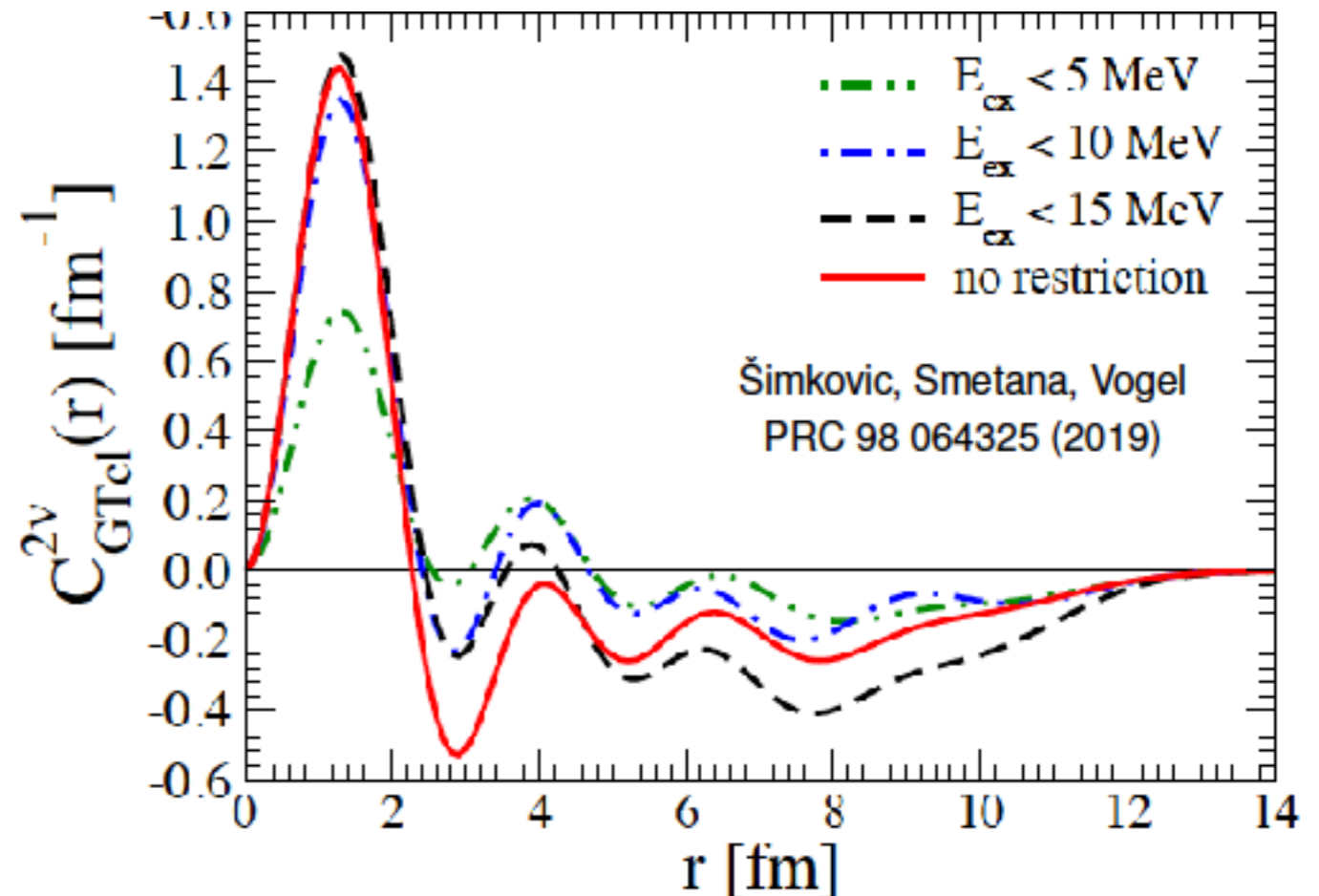
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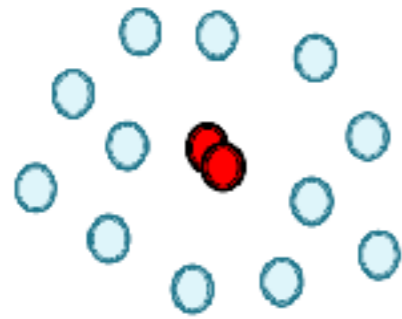
Reduction (QRPA) or not (SM, etc.) of DGT with respect to  $E_{ex}$  translates to very different radial dependence



# Wave function factorization



# Wave function factorization



$\mathbf{r}_{12} \rightarrow 0$

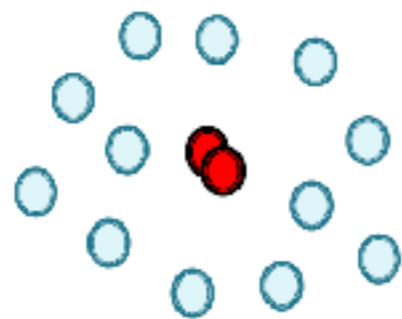
$$\Psi_n^A(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A) \sim \phi(\mathbf{r}_{12}) \chi_n^A(\mathbf{R}_{12}, \mathbf{r}_3, \dots, \mathbf{r}_A)$$

Universal  
2-body physics

state-dependent  
many-body physics



# Wave function factorization



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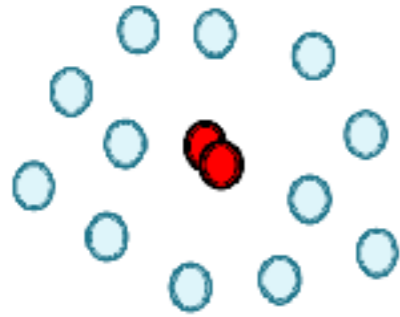
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Consider **high**  $q$  ( $\gg \Lambda$ ) components of **low**  $E$  ( $\ll \Lambda^2$ )  $A=2$  wf's

$$Q|\psi_n\rangle = \frac{1}{E_n - QH} QVP|\psi_n\rangle$$

$$\psi_n(q) = \int_{\Lambda}^{\infty} d^3q \int_0^{\Lambda} d^3p G_Q(q, q'; E_n) V(q', p) \psi_n(p)$$

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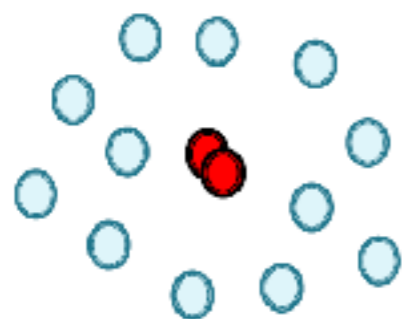
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scale separation:  $G_Q(q, q'; E_n) \approx G_Q(q, q'; 0) + \dots \quad V(q', p) \approx V(q', 0) + \dots$

# Wave function factorization



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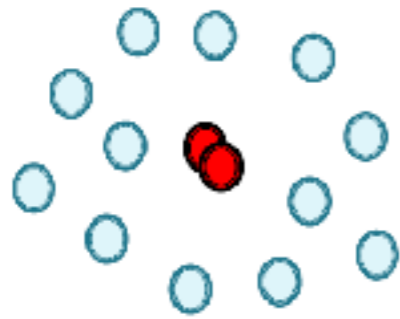
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scale separation:  $\Rightarrow \quad \psi_n(q) \approx \gamma(q; \Lambda) \int_0^{\Lambda} d^3p \psi_n(p) + \eta(q; \Lambda) \int_0^{\Lambda} d^3p p^2 \psi_n(p) + \dots$

# Wave function factorization



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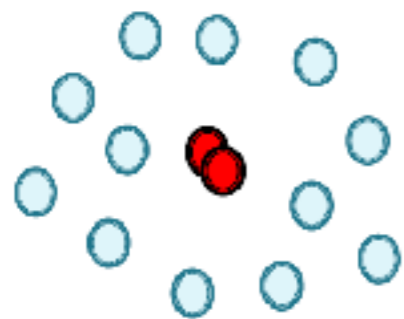
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**Operator Product Expansion**

Wilson coefficients

ME's of (smeared) local operators



$$\mathbf{r}_{12} \rightarrow 0$$

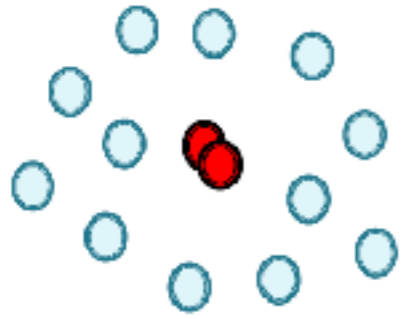
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Cruz-Torres et al. arXiv:1907.03658

E.g., pair density

$$\rho_{2b}(\mathbf{r}) = \langle \Psi | \sum_{i < j} \delta(\mathbf{r} - \mathbf{r}_{ij}) | \Psi \rangle$$

# Wave function factorization



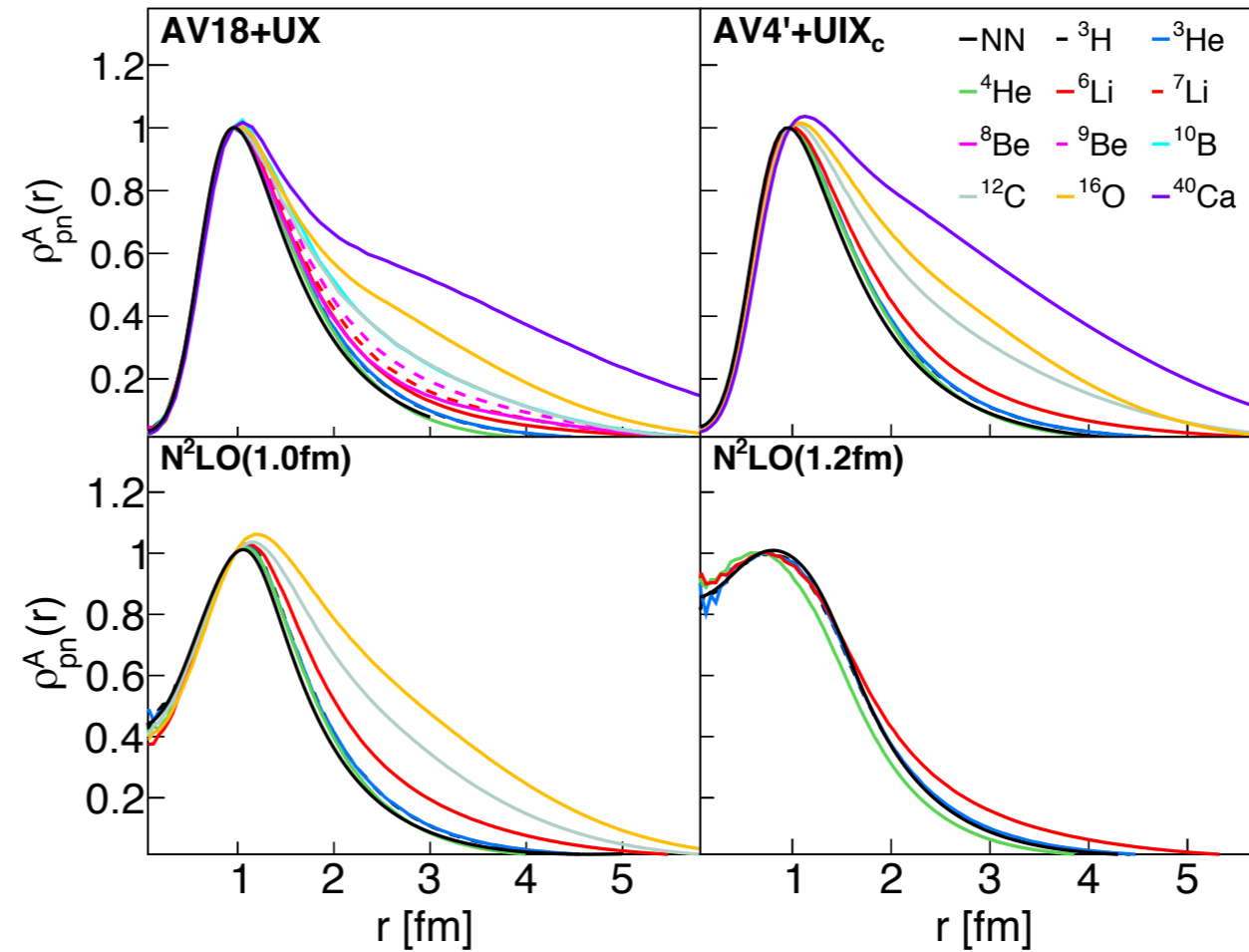
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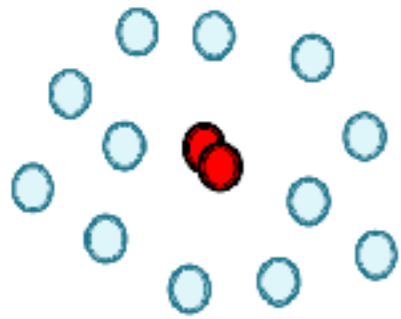
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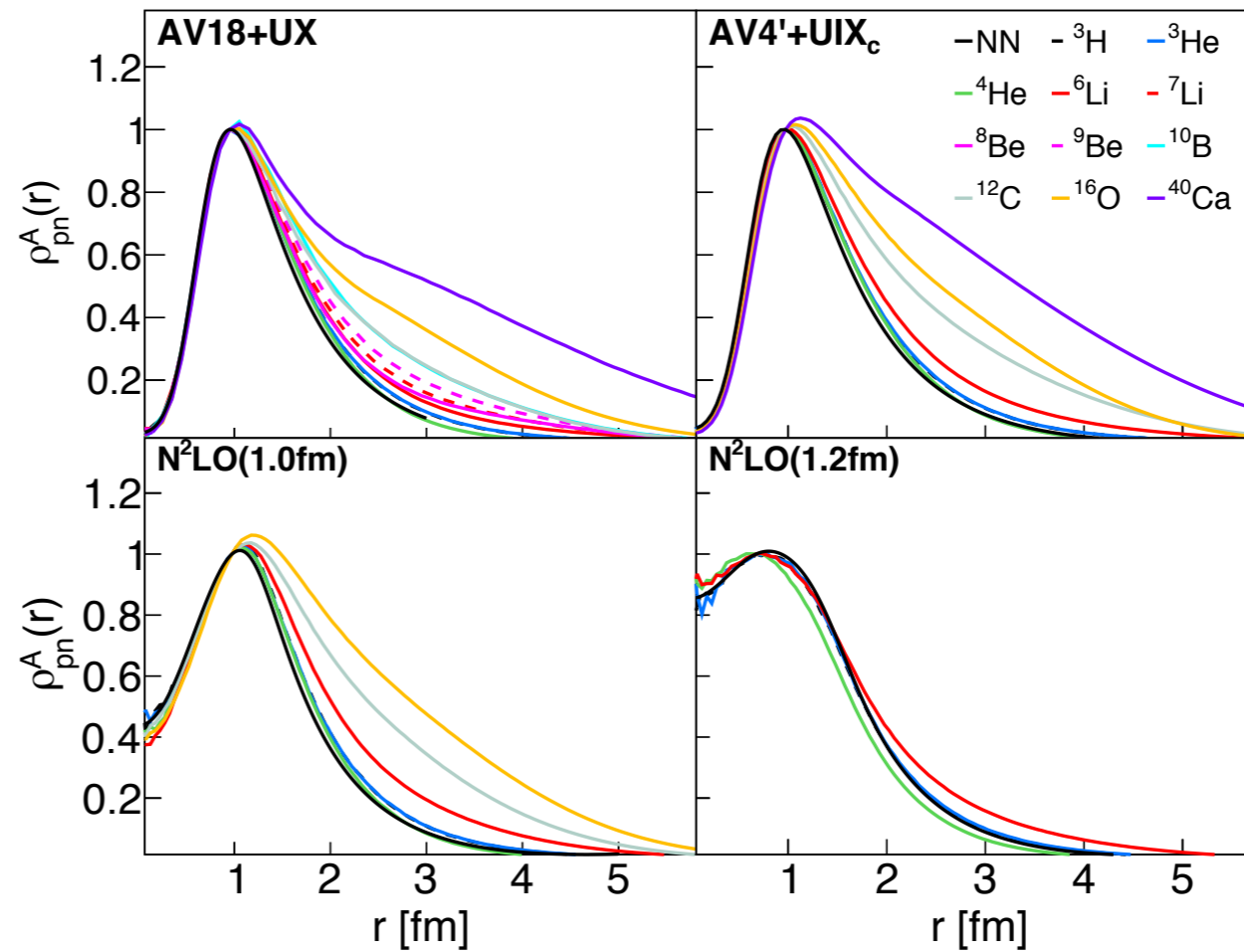
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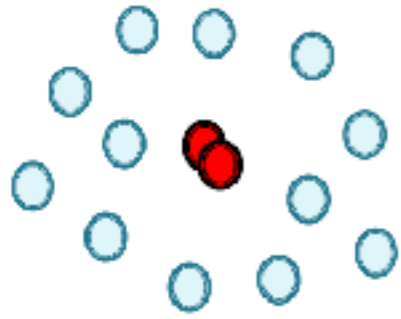
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Universal short-distance shape from wilson coeff  $\phi(r)$

A-dependent scale factor controlled by  $\chi$

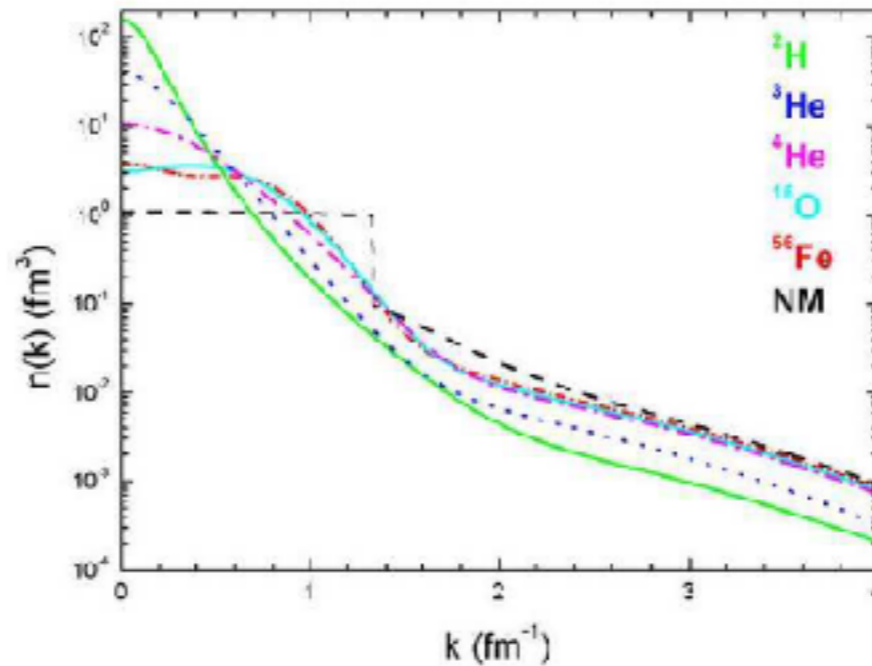
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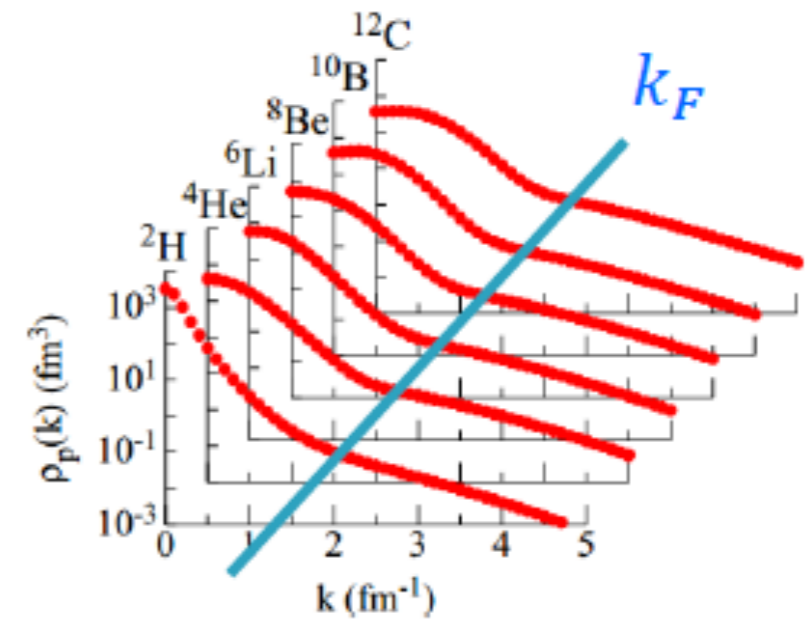
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e.g., 1- and 2-body momentum distributions



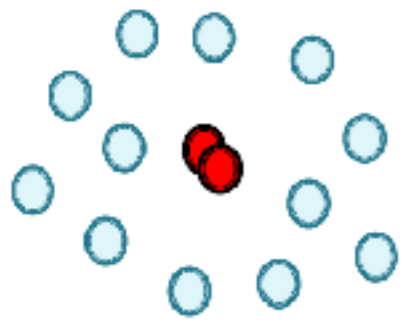
[From C. Ciofi degli Atti and S. Simula]



R.B Wiringa et. al., Phys. Rev. C 89, 024305 (2014)



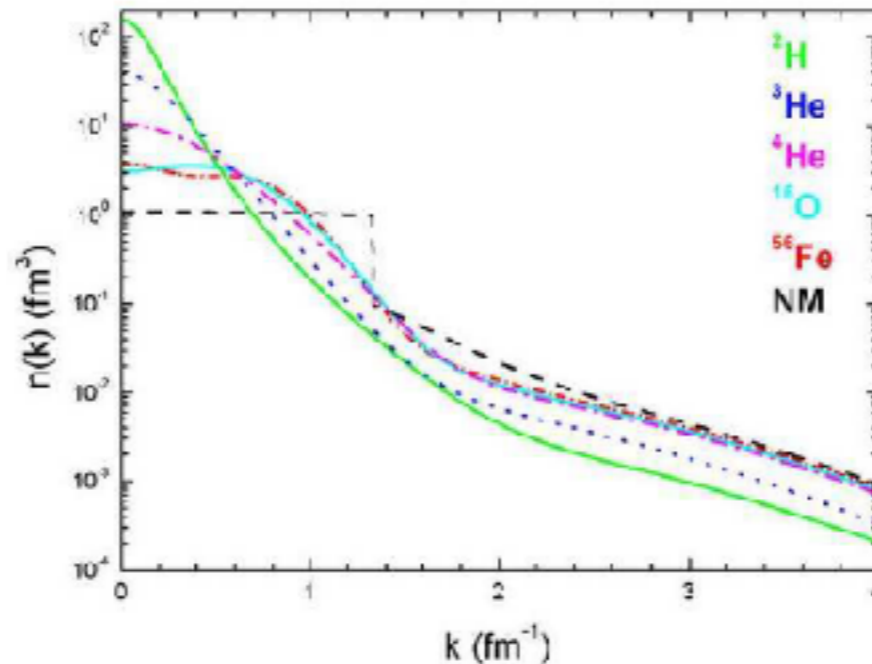
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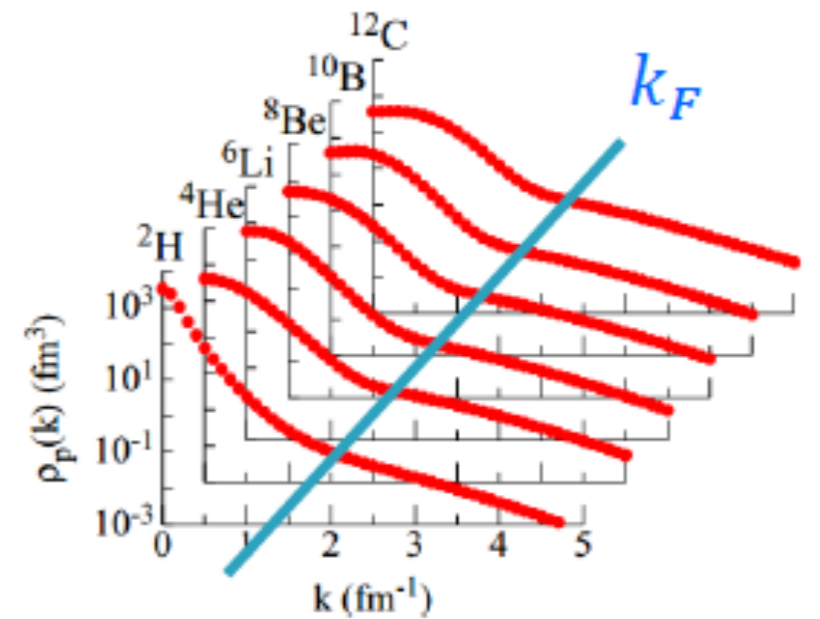
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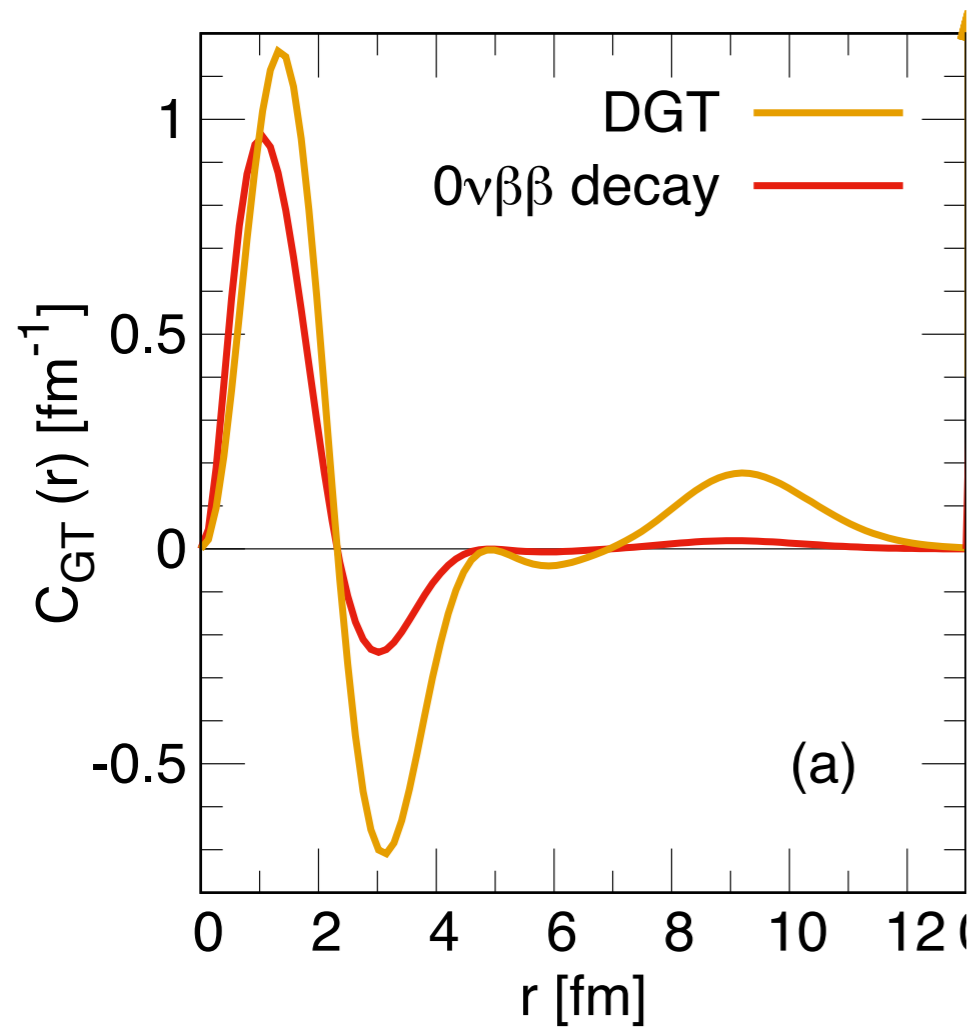


R.B Wiringa et. al., Phys. Rev. C 89, 024305 (2014)

Universal high-q tails, controlled by wilson coeff  $\phi(\mathbf{r})$

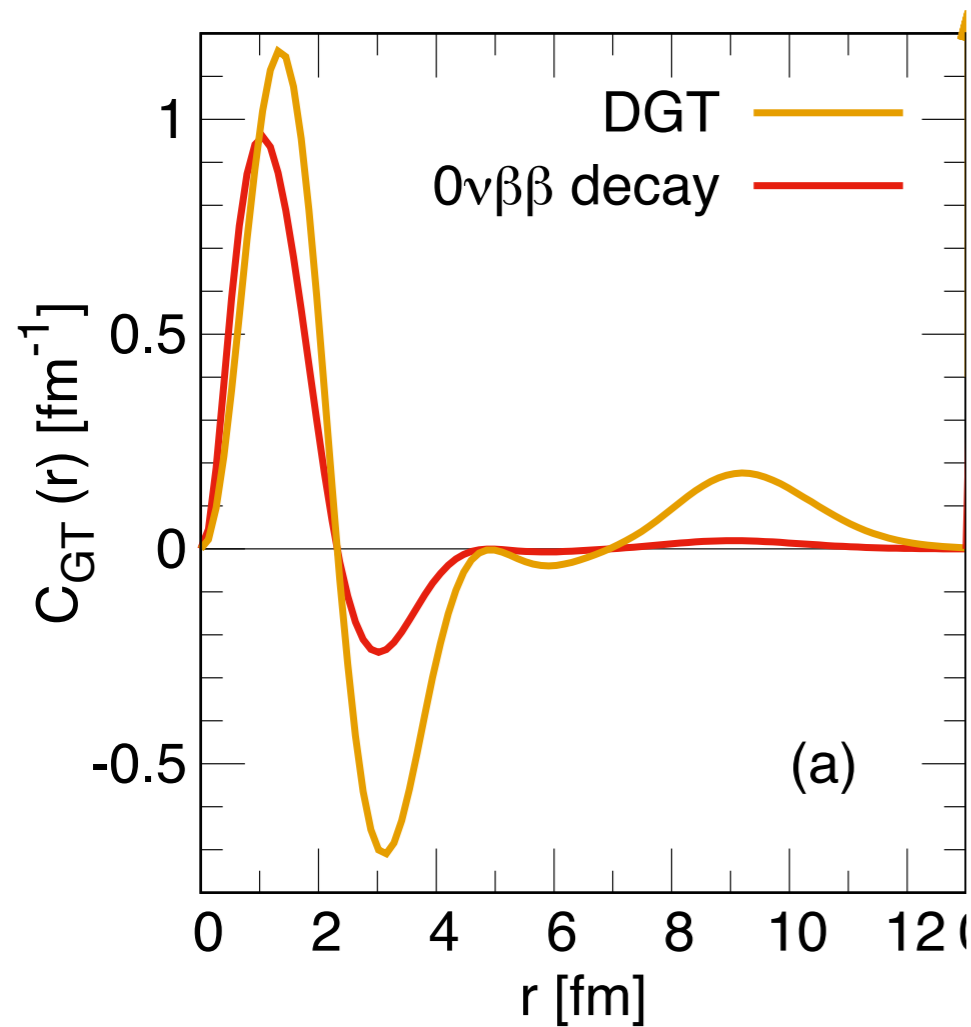
A-dependent (q-independent) scale factor controlled by  $\chi$

# Connection with $0\nu\beta\beta$ -DGT correlation?



$$\mathcal{M}_{0\nu,GT} = \frac{N(N-1)}{2} \int d1 \cdots dA \Psi_f^\dagger(1,2,\dots,A) H_{GT}(r_{12}) \sigma_1 \sigma_2 \tau_1 \tau_2 \Psi_i(1,2,\dots,A)$$

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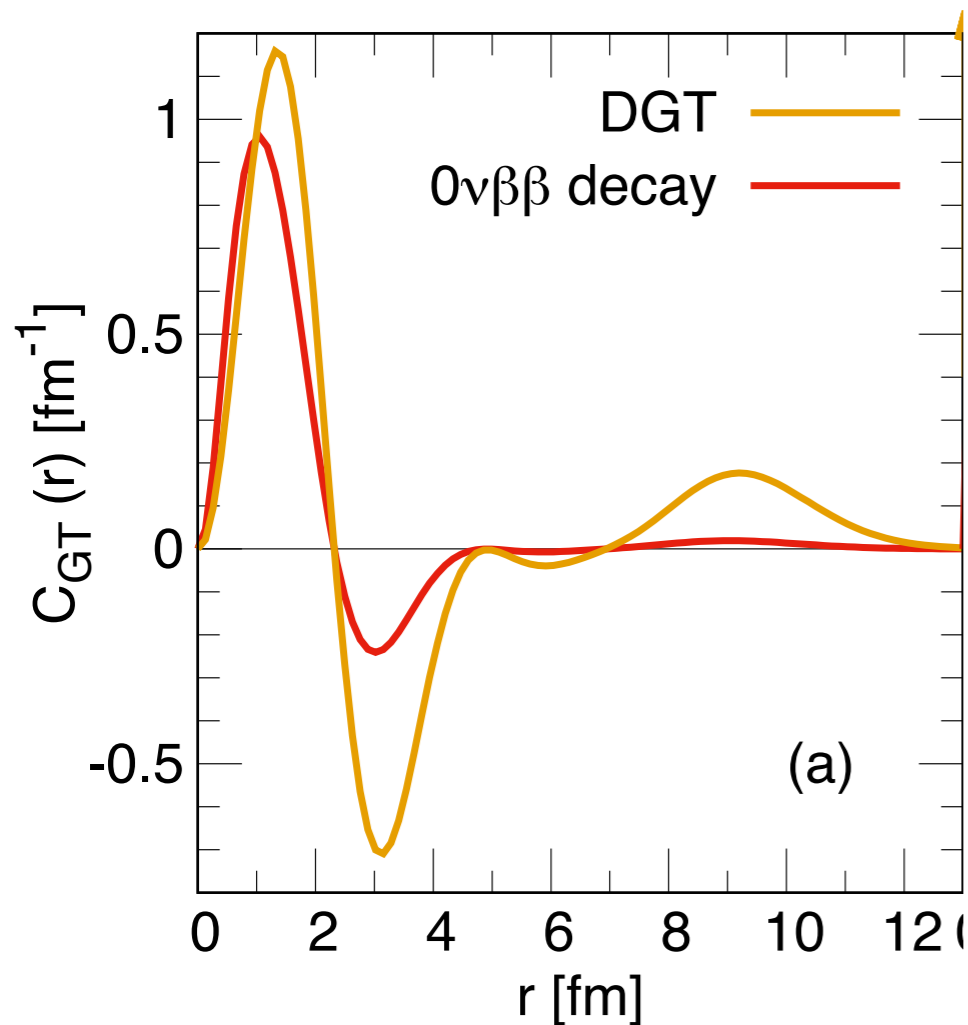
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$$\Psi_i(1,2,\dots,A) \sim \phi_{nn}(r) \chi_i(R; 3,\dots,A)$$

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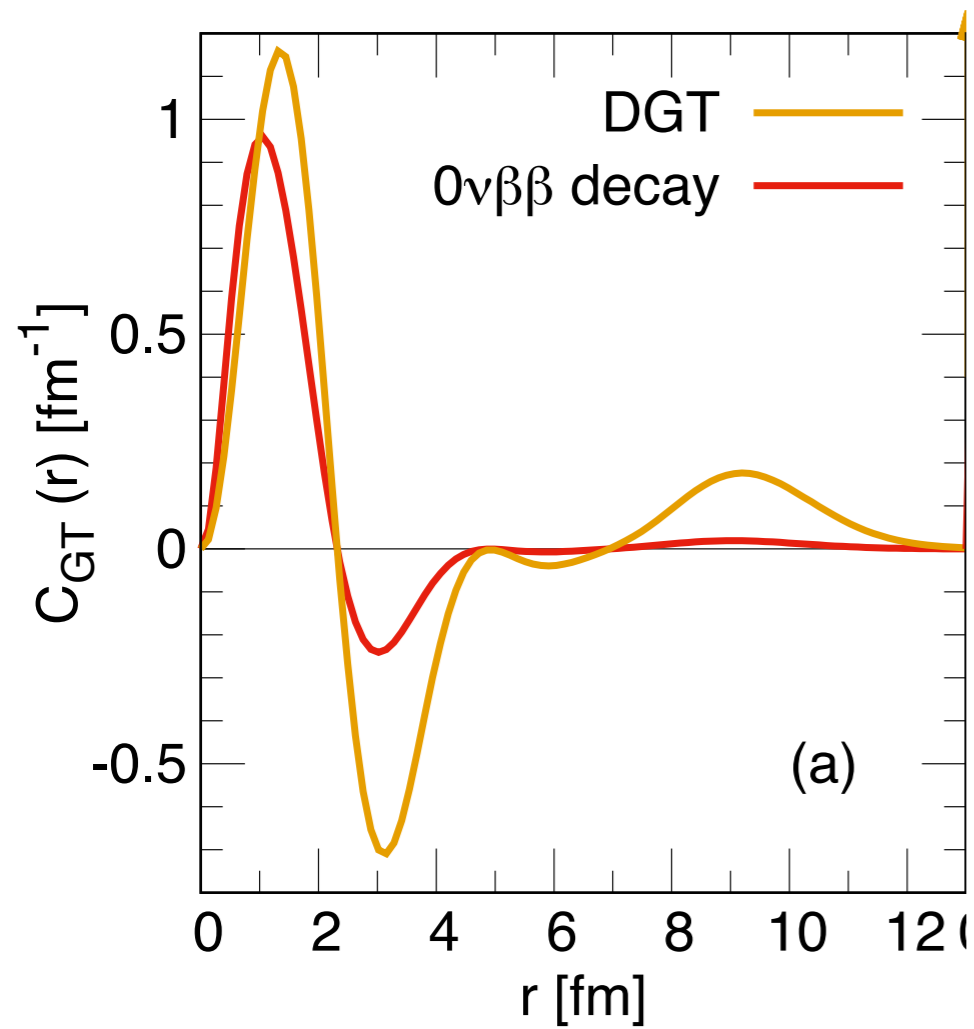
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$$\frac{\mathcal{M}_{0\nu}}{\mathcal{M}_{DGT}} \sim \frac{\int^{r_c} |\phi(r)|^2 H_{GT}(r)}{\int^{r_c} |\phi(r)|^2}$$

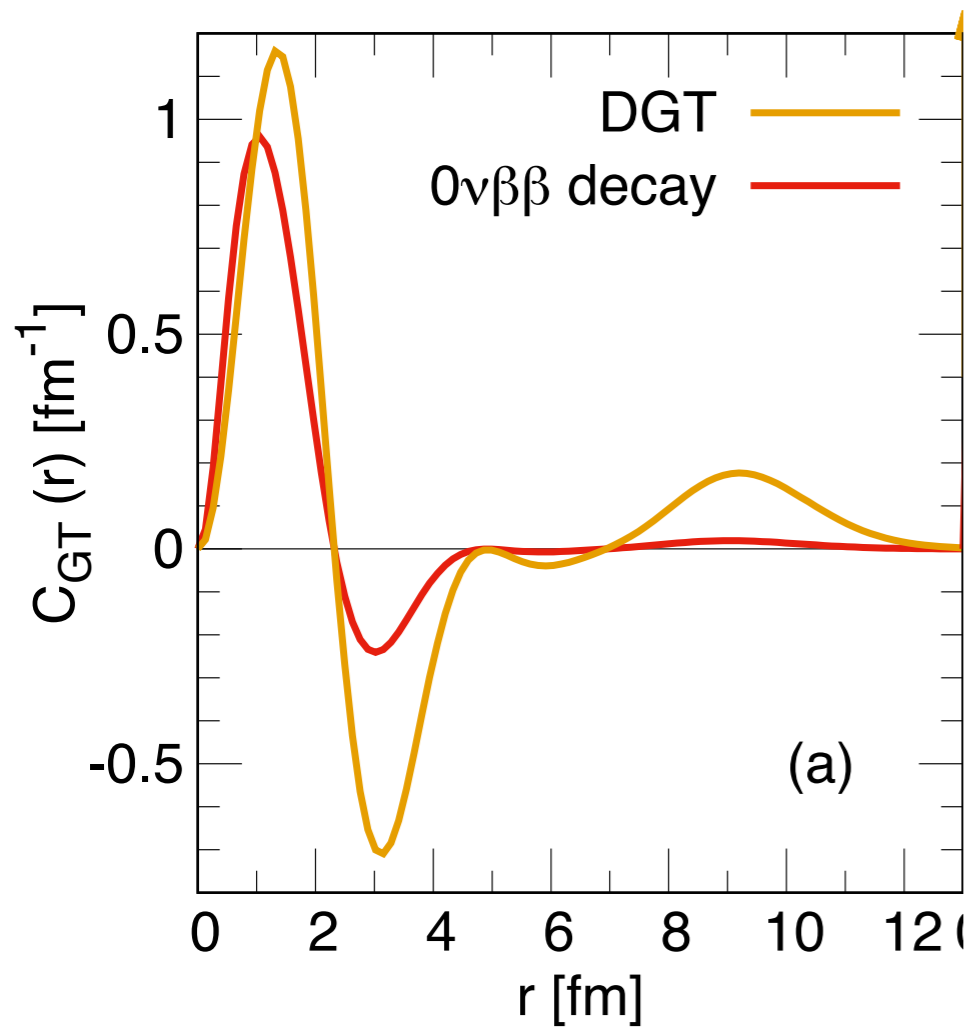
linear correlation follows  
many-body physics cancels in ratio

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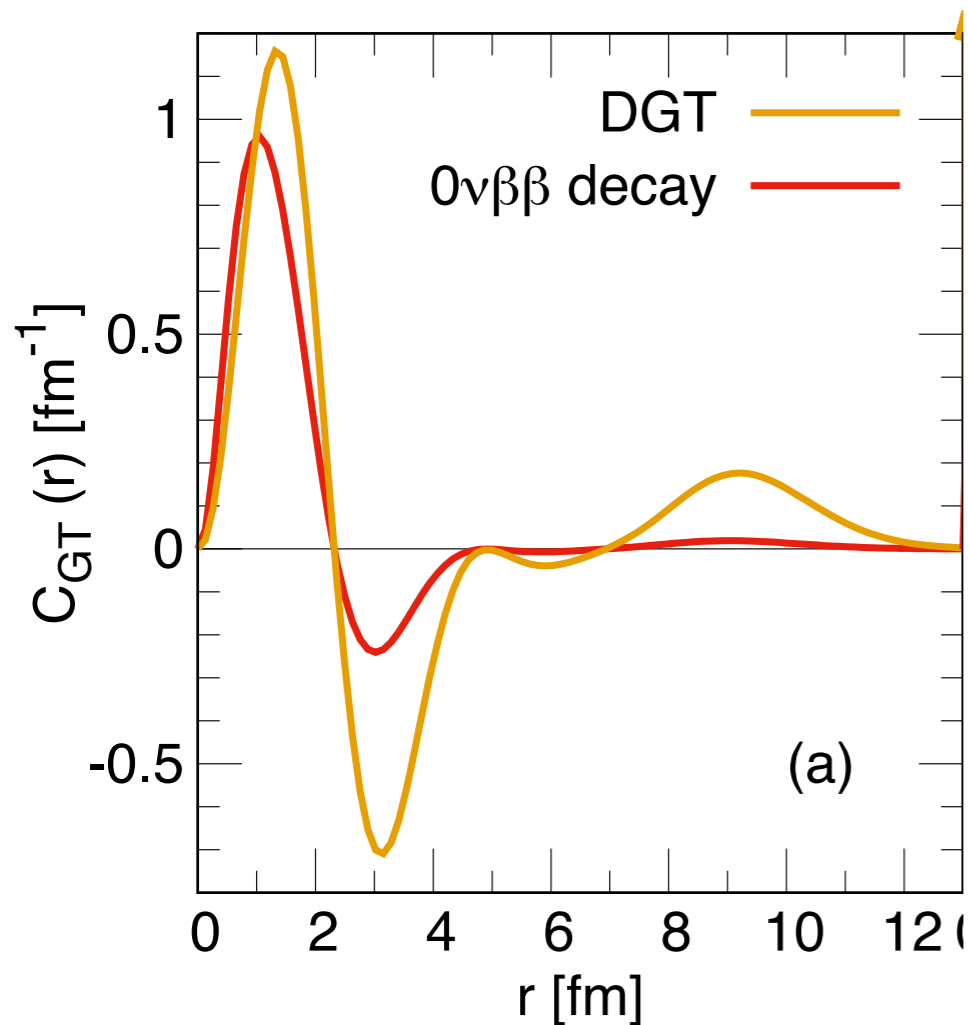
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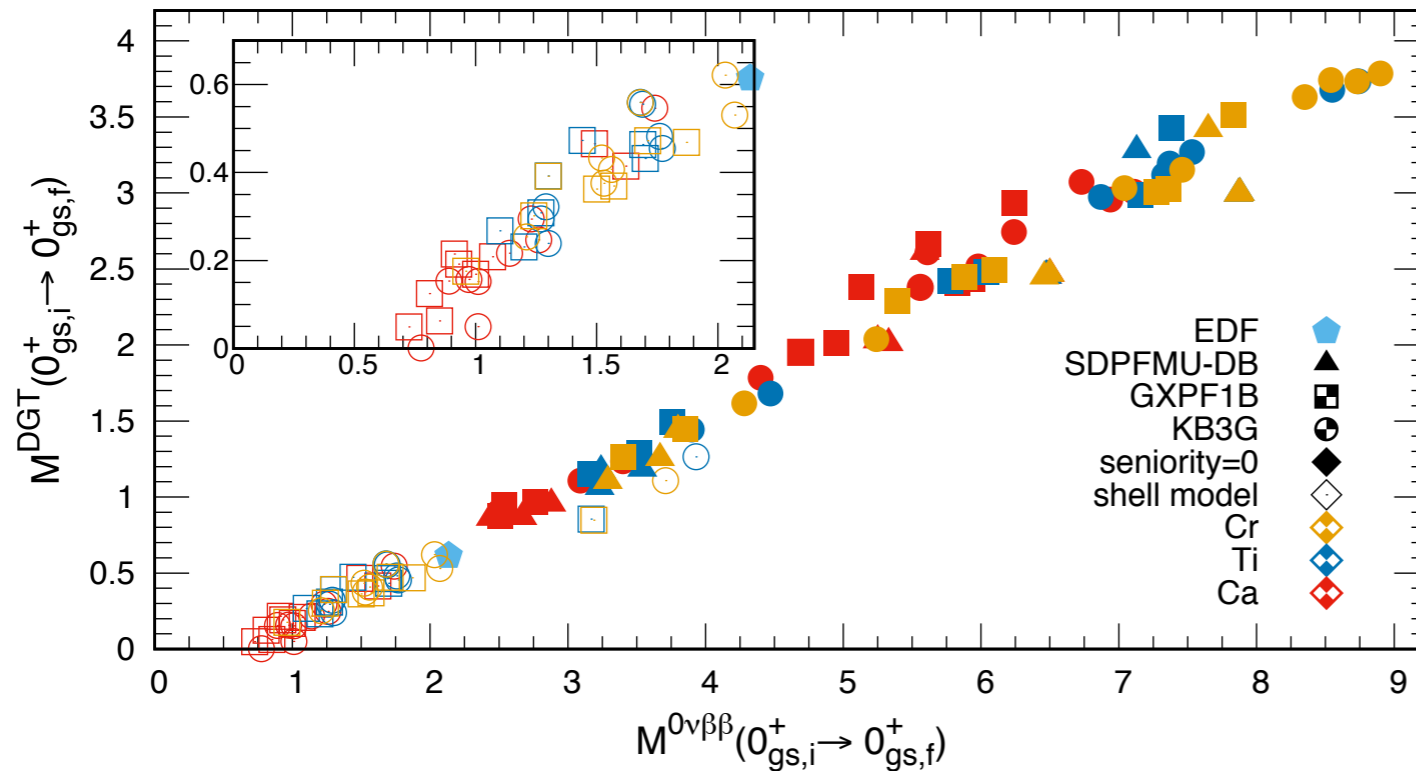
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# Simple estimates from factorization



$$\frac{\mathcal{M}_{0\nu}}{\mathcal{M}_{DGT}} \sim \frac{\int^{r_c} |\phi(r)|^2 H_{GT}(r)}{\int^{r_c} |\phi(r)|^2}$$

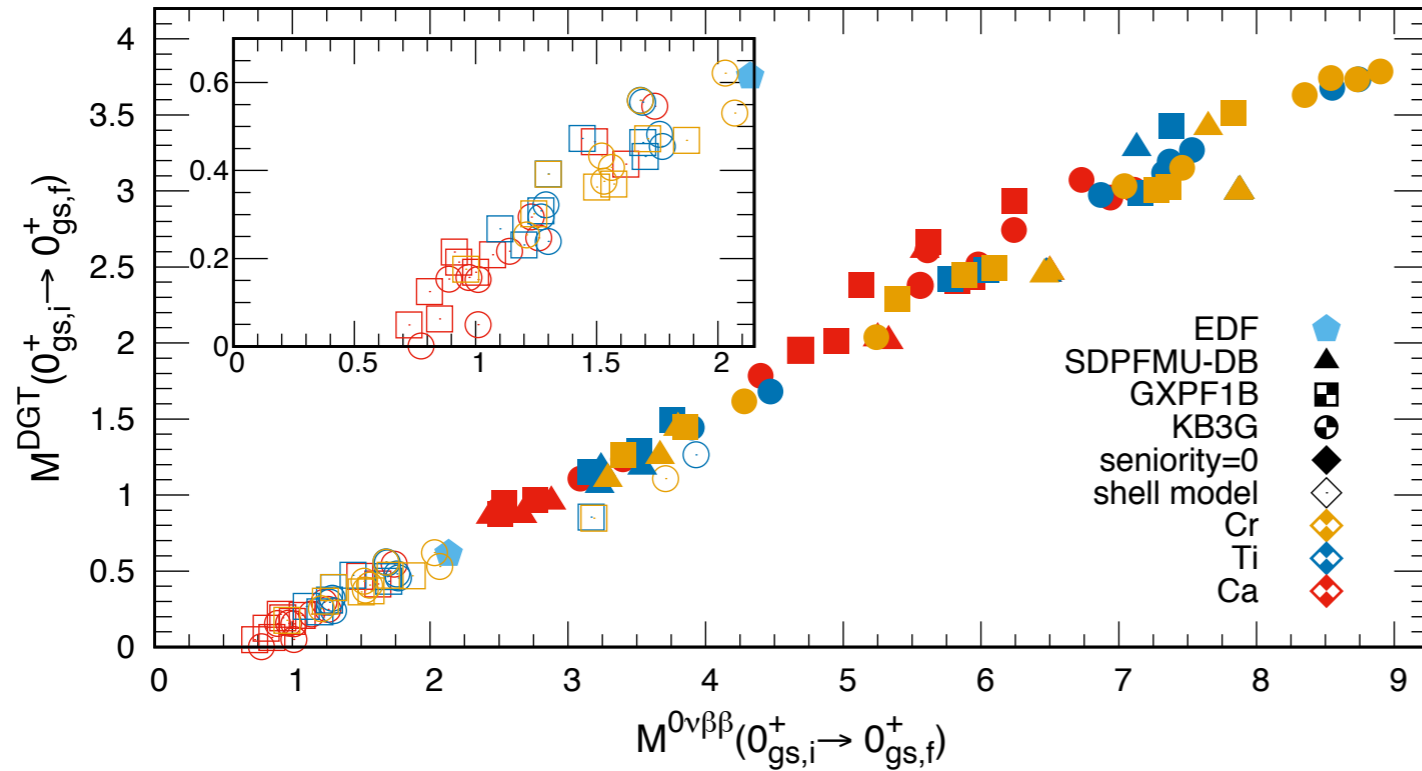
Universal  $\Phi$  formally given by  $E=0$  2-body w.f. (c.f. Weiss et al. PRC 92 (2015))

Scale and scheme-dependent

Shell model, EDF, IBM, etc. are intrinsically “low resolution” pictures. Try some simple s-wave  $\Phi$ 's with low-resolution scales



# Simple estimates from factorization



$$\frac{\mathcal{M}_{0\nu}}{\mathcal{M}_{DGT}} \sim \frac{\int^{r_c} |\phi(r)|^2 H_{GT}(r)}{\int^{r_c} |\phi(r)|^2}$$

observed slope  $\sim .59$

First attempt:

$\Lambda$ MeV	$\frac{1 - e^{-r\Lambda}}{r}$	$j_0(\Lambda r)$
200	0.47	0.44
100	0.51	0.53
50	0.53	0.55
20	0.54	0.56
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# Simple estimates from factorization

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Calculations missed factor of  $g_A^2 = 1.59$ . Numbers now span between  $.28 - .36$

- Factorization/OPE ideas explains(?) (or is at least consistent with) the linear  $0\nu\beta\beta$ -DGT correlation
- First leading order OPE estimates were shockingly good (5-20% level) , until the damn extraneous  $g_A^2$  factor was noticed.
- Still, the trend is not horrible. Slopes  $\sim .28 - .36$  using “reasonable” universal  $\Phi$  appropriate to low-resolution SM/EDF/IBM
- Can we improve this, e.g., subleading OPE terms  $\Rightarrow$  deviations from linear scaling?
- Do the ab-initio (IMSRG/IM-GCM/CC) results likewise scale with DGT?

# Wave function factorization

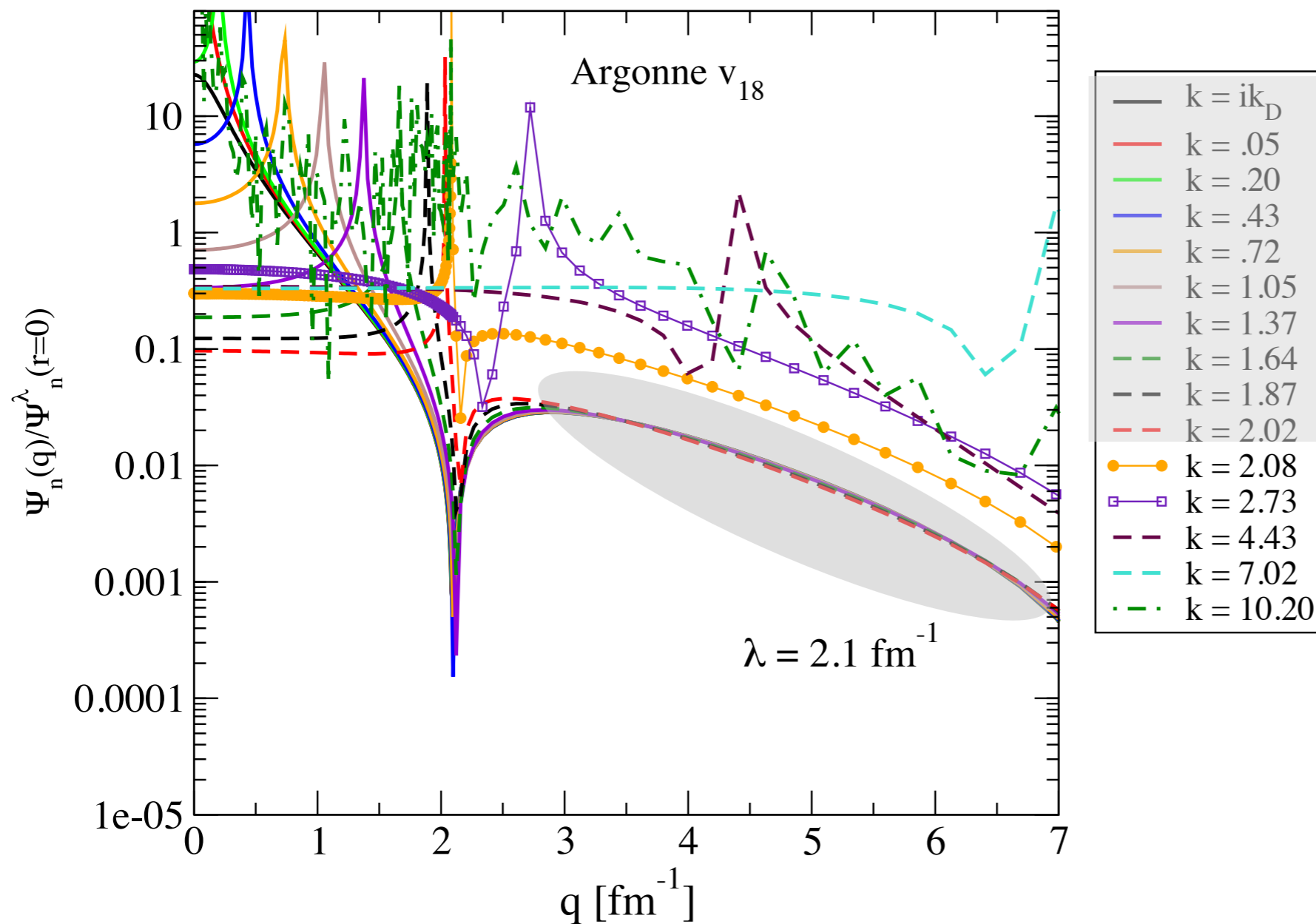
LO:  $\psi_{\alpha}^{\Lambda_0}(\mathbf{q}) \approx \gamma(\mathbf{q}; \Lambda) \int_0^{\Lambda} d^3p Z_{\Lambda} \psi_{\alpha}^{\Lambda}(\mathbf{p})$   state-independent ratio  
for well-separated scales

$$\frac{\psi_{\alpha}^{\Lambda_0}(\mathbf{q})}{\psi_{\alpha}^{\Lambda}(\mathbf{r} = 0)} \sim \gamma(\mathbf{q}; \Lambda)$$

$$|E_{\alpha}| \lesssim \Lambda^2 \quad |\mathbf{q}| \gtrsim \Lambda$$

# Wave function factorization

LO:  $\psi_{\alpha}^{\Lambda_0}(\mathbf{q}) \approx \gamma(\mathbf{q}; \Lambda) \int_0^{\Lambda} d^3p Z_{\Lambda} \psi_{\alpha}^{\Lambda}(\mathbf{p})$   $\rightarrow$  state-independent ratio for well-separated scales

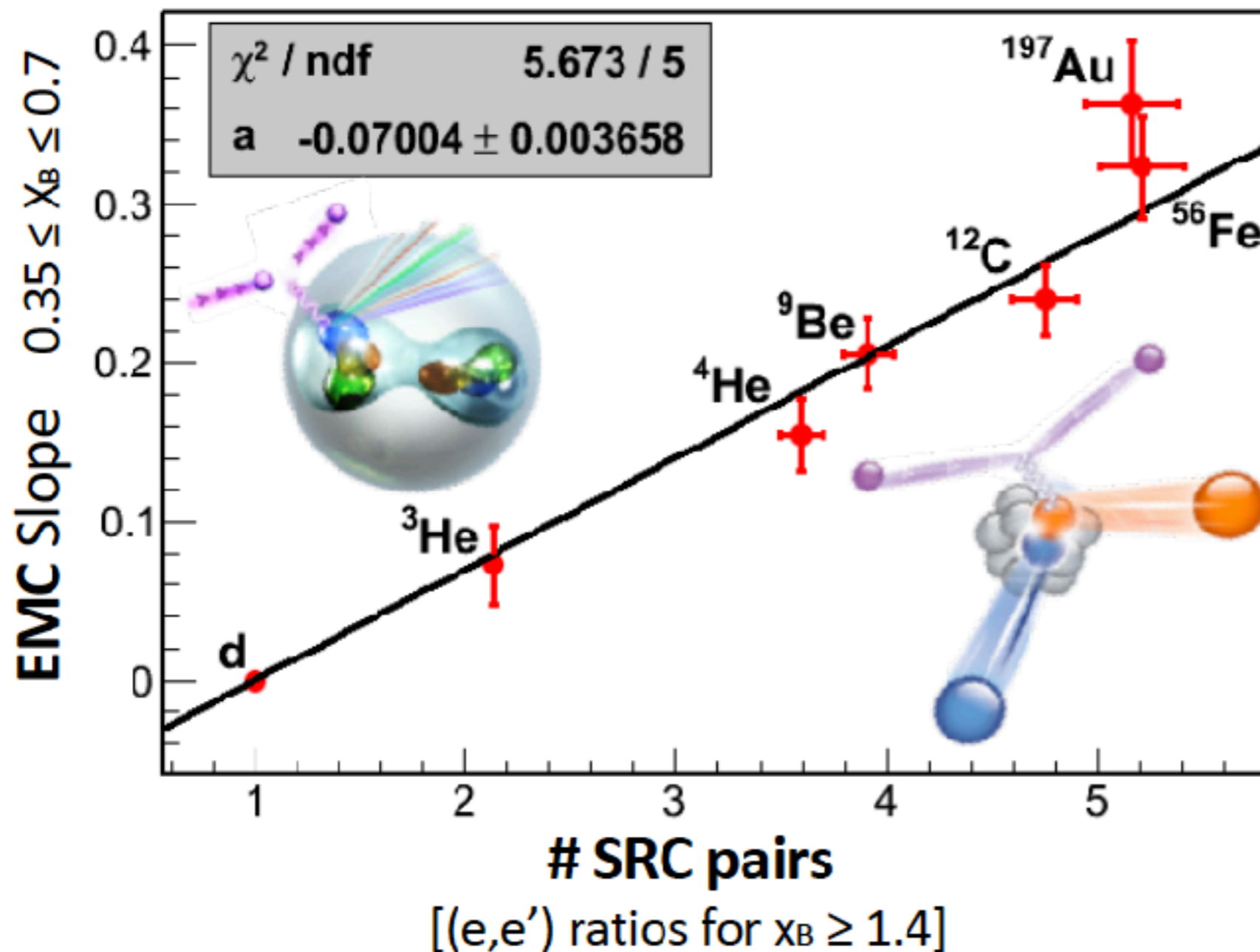


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# Empirical correlation of EMC effect

Hen et al., RMP (2017); Hen et al., IJMPE (2013); Hen et al., PRC (2012);  
Weinstein, Piassetzky, Higinbotham, Gomez, Hen, and Shneor, PRL (2011).



Why should 2 seemingly unrelated processes be linearly related?